On the preparation of a prototype VCN source

E. Lychagin







The prototype should demonstrate the properties and characteristics of both the converter and the reflector.

The main characteristic of the converter is the dependence of VCN production rate on the energy of the VCN.

The main characteristic of the reflector is the gain in the VCN flux in limited solid angle along the axe and its depending on the VCN energy.

The experiment parameters:

Neutron beam	PF1b, ILL		
VCN converter	o-SD ₂ , 5K, ~4 cm³		
Estimated production rate Φ_{VCN}	1.6·10 ⁶ n/s (v<100 m/s); 2 ·10 ⁵ n/s (v<50 m/s).		
Estimated detection rate without reflector Φ_{det}	16 n/s (v<100 m/s); 8 n/s (v<50 m/s).		
Estimated gain factor	3÷6 (v<100 m/s); 4÷15 (v<50 m/s).		



Main elements of the experimental set-up:

- Beam collimation system (standard equipment PF1b);
- Cryostat for VCN converter and reflector (under developing in JINR);
- Ortho-para converter for D₂ (under developing in JINR);
- Powder for reflector is producing by PTI and UCA (will be in reports A.Vul' and M.Dubois today)
- VCN chopper (was developed in JINR and located in ILL);
- VCN detector (standard equipment of ILL).





Schema of cryostats for VCN converter and optho-para converter

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The design work was started but temporarily stopped till a final decision have been done concerning the design of the central part where the converter and reflector are located.



Central part





The teflon tube with 1 cm internal diameter and wall thickness 1 mm was chosen as deuterium container. This tube will have cold aluminum end to freeze the gas.

Internal wall of volume for powder reflector will prodused by teflon with 200 mu thickness.

Loss fac	tor	P _{loss} (V _{VCN} = 100 m/s)	P _{loss} (V _{VCN} = 50 m/s)	
Teflon:	P _{loss} (CF ₂)	3·10 ⁻³	6·10 ⁻³	
SD ₂ :	P _{loss} (D ₂)	1.4·10 ⁻³	2.8·10 ⁻³	
Holes:	P _{loss} (holes)	1.9·10 ⁻²	1.9·10 ⁻²	
Powder of diamond:	P _{loss} (DNDs)	~0.15	~4.10-2	
TOTAL:	P _{loss}	~0.17	~6.8.10-2	

Material	density, g/sm ³	м	σ _a , barn (2200 m/s)	molecular density n *10 ²³ , 1/cm ³	n* σ_a (50 m/s), 1/cm	thickness to suppress flux (50 m/s) at 10%, cm	Wall thickness at 50 reflections, um
Teflon	2,200	50,0	0,023	0,265	0,026	3,98	265,4965775
Mg	1,738	24,3	0,063	0,431	0,119	0,88	58,85094118
Ве	1,848	9,0	0,008	1,236	0,041	2,55	169,9277907



Raman measurements



We made a gas test cell and started the game with measuring ща Raman spectra to determine ortho to para state the ratio in the deuterium.





Nearest plans (for next 6 month):

- Development of working drawings of the central part.
- Test of ortho-para conversion at 20K (et test refrigerator).
- Adjustment of the complete installation scheme and finalized list of needed equipment.
- Completion of the development of working drawings for cryostats, supports, neutron shielding.
- Start manufacturing of the set-up elements.



On the development of a soliddeuterium VCN convertor

Ekaterina Korobkina

North Carolina State University, Raleigh, USA





Nuclear spins orientation: Ortho-deuterium

- D2 molecules exist in two isomeric states para-D2 (spin J=1) and ortho-D2 (spin J=0)
- During condensation the rate of conversion to the equilibrium state is extremely slow
- Neutrons can introduce such conversion
- As result, VCN/UCN gain energy of several meVs and are lost
- Therefore we need to use pre-converted orthodeuterium to growing sD2 crystals



FIG. 1. UCN upscattering cross section vs temperature of solid D_2 . The one-phonon annihilation cross section in an ortho- D_2 solid (solid curve) and in a para- D_2 solid (dashed curve) are plotted. The dashed-dagger line is the temperature independent UCN upscattering cross section involving $J=1\rightarrow 0$ relaxation not coupled to phonons in a para- D_2 solid.





Solid deuterium grows modes

- Freezing from the liquid phase (LANL, PSI)
- Sublimation from the vapor phase at T_s:
 - at the operational temperature: T_s=T_{op} (Mainz Uni)
 - above operational temperature: T_s > T_{op}



• for UCN/VCN source $T_{op} = 5K$

Solid deuterium properties

- The SD2 crystal is a quantum crystal with a large zero T energy
- It demonstrates thermally activated diffusion above 8K (D(D2)=7x10⁻⁴e^{-290/T}) as well as surface diffusion
- It demonstrates significant shrinkage from 18.3K to 12K

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 Crystals grown from vapor can be a mixture of hcp and fcc lattices



Solid deuterium properties

 SD2 exhibits so called "triple point wetting " (or Stranski–Krastanov film growth), i.e. it wets surface completely above the triple point, while below only films of several monolayers can be grown and the rest grows as crystallite islands



Figure 7. Wetting curves of D_2 on Au are shown. A similar behaviour has been observed for other kinds of adsorbates. In the inset, the same curves are plotted on a logarithmic scale.



Figure 1. Cross-section views of the three primary modes of thin-film growth ^Δ including (a) Volmer–Weber (VW: island formation), (b) Frank–van der Merwe (FM: layer-by-layer), and (c) Stranski–Krastanov (SK: layer-plus-island). Each mode is shown for several different amounts of surface coverage, Θ.





Solid deuterium geometry for the VCN prototype

- The SD2 crystal has to be grown horizontally
- dimensions I.D.1cm
 x 5 cm long





NCSU study of growing SD2 crystals

- Visual observation
- Temperature sensors inside and outside the cryocontainer which allowed us to reconstruct temperature profile of the walls





Heater wrapped at the top of container



T-sensor locations at the cry-container plus a sensor on the He return line



Growing SD2 from vapor at UCN source operational temperature





Simulation of Mainz UCN source

- this run was to simulated Mainz UCN source condensation of SD2 with cold (6K) bottom of container and slow D2 flow rate
- 5.4K(He inlet)/ 8.2K(Top)/ 7.2K(He outlet):
 - Small flow (0.3 l/m) produced dense multicrystall, optically opaque
 - Higher D2 flow >1
 I/m produced snowflake-like mass
- when T of container was increased (5.4K/ 11.2K/ 8.3K) and crystal annealed, D2 flow =>0.8 I/m results in visibly shiny surface





After two days of annealing at 12K



2018, melting-refreezing, resulting in grainy crystal

• this melting-refreezing was done faster than 2016 run and the crystal was cooled down to 9.5K for annealing







Growing from vapor above the UCN source operational temperature: case 1





Condensation at T close to T_{triple} and annealing at 11.5K

- Mar 15 condensation, no heaters, small cooling power, 9/ 17.5/ 17.5, D2 flow 0.8 l/m:
 - ideal transparent
- Macnystad-18, evolution to a blob at 11.5K:
 - amazing mobility and tendency for avoiding warmer surfaces
 - Mar 16 condensation, with heaters on, 8.5/18/ 15K, 1 l/m



- Conclusions:
 - optically clear crystal
 - high mobility and tendency to pull off the foreign objects

Mar 22- Mar 25 - condensation/conditioning



Mar 22 -25 condensation, with heaters, 7.75/17/ 18.5-19, flow 1 l/m

- transparent crystal
- all probe sensors covered, about 1050 cc total

• Conclusions:

- it is possible to grow optically good crystal from vapor at warmer temperatures and flow 0.8-1 l/m
- need to be concerned about crystal shape evolution with time





ISINN 2017

Conclusion from first two condensation



- at T close to T_{triple} , :
 - crystal is clear
 - SD2 move and re-shaping itself at 12K due to high diffusion
 - need to be concerned about crystal shape evolution with time

- at T<6K SD2 does not move
- but the growth is very inhomogeneous due to "triple-point-wetting"
- annealing at 12K improved crystal bulk





Growing from vapor above the UCN source operational temperature: case 2





Growing from vapor above operational T: SD2 shapes and quality vs Temperature profile of the cryo-container



Growing from a vapor at intermediate T

- April 11-13 condensation, 8.3/ 16.2/13, 0.8 l/m
- attempt to reproduce Mar 15 condensation (9/ 17.5/ 17.5)
- it was unsuccessful, because D-return T was nor taken into account and the container was too cold
- very interesting run for discussion of crystal properties vs T and crystal shapes

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How we can extract crystal shape from our data

• Position of the T- probe





How we can extract crystal shape from our data

• Hight at probe: coverage of probe sensors



How we can extract crystal shape from our data

- We know total Volume
- We can estimate outer diameter of the crystal





Shape reconstruction

Mar 15, Annealed overnight at 12K. High mobility





Shape reconstruction (about 350cc)

March 22















CREMLIN+ WP3, 2021

Shape reconstruction (1050cc vs 964cc)

March 25











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Conclusion from SD2 shape reconstruction

The only difference between March and Apr 11 runs was temperature distribution of the container bottom



Crystal shapes and transparency are extremely temperature dependent!



TRIUMF seminar, 2019

How to simulate T-distribution of container walls

- T-top and T -inlet are fixed on surfaces shown on the left
- Then temperature of the bottom circle is raised above D-inlet until 1 cm reading matches experimental value





TRIUMF seminar, 2019

Mar 22 condensation, 1cm=16K; 7.75/17/18.5-19K; from simulations bottom centre =14K;

 1cm temperature can be reproduced only when assuming that bottom circle is at least 14K









April 11 simulation, D2 flow on and off

- 1cm temperature can be reproduced only when assuming that bottom circle is at 9.3K
- Transparent region start from above 12K
- after D2 flow was off, 1 cm dropped to 11K, while the T of container has not changed, it consistent with Bottom circle at the same T as D-inlet





CREMLIN+ WP3, 2021
April 25



- 9.5/18/
 16.2,1cm=14.5K, about 340cc total
- 1cm temperature can be reproduced only when assuming that bottom circle is at 12K
- **Conclusion**: to grow transparent crystal, cryostat walls needs to be above 12K





ISINN 2017

Overview of 2016 results

Summary/conclusions from 2016

- There is a significant difference in crystal quality vs D2 flow, depending on cryostat temperature:
 - flow > 0.8 l/m & cryostat at 6K result is snowflakes;
 - cryostat above 10K good shiny crystal
- Quality of crystal depends critically on the growth temperature:
 - growing at T =6K results in a disconnected polycrystalline structure
 - growing or annealing crystal at T =9.5K results in a one solid while very grainy structure
 - for good crystal the coldest spot should be at T > 12K
- We observed extremely high mobility of SD2 even at 11.5K (couple of mm per hour):
 - large temperature gradients result in deformed crystal shapes
 - better to avoid use of heaters



Growing perfect sD2 from vapor above the UCN source operational temperature following 2016 conclusions in 2018





May condensation

- this run was done in one step, with one continuously decreasing flow of D2
- initial Ts: 8K/ 18K/ 18K; no heater on during condensation; Ts controlled by cold He-flow tuning.
- final Ts: 7K/ 15K/ 15K, D2 probes at 14K gradient only 1K



May condensation

- this run was done in one step, with one continuously decreasing flow of D2
- initial Ts: 8K/ 18K/ 18K; no heater on during condensation; Ts controlled by cold He-flow tuning.
- final Ts: 7K/ 15K/ 15K, D2 probes at 14K gradient only 1K





May condensation: cooling





May condensation: accidental annealing

- I set up heat shooting, as result at one point there was a bible in LHe line and the crystal was annealed to 14K without significant increase of pressure
- All surface facets were gone after such annealing





quality and shape of the crystal can be predicted if we know the Temperature distribution of the crycontainer walls



Backup slides





Growing from liquid phase (melting and re-freezing)





SD2 growing test 2016: melting



D2 pressure 43 mbar



SD2 growing test 1: start of freezing









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SD2 SD2 freezing after melting



SDZ freezing after meiting, accidental warm



SD2 after annealing overnight at 16K



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SD2 after cooling down to 5K





- It seems that because of annealing we were able to avoid contraction strain and obtain crystal of the quite good quality
- during annealing crystal was re-shaping and moving from the walls







On simulation of VCN sources

L. Zanini on behalf of

Zs. Kókai, J.I. Márquez Damián, V. Santoro, L. Zanini

This project has received funding fram the European Union's Horizon 2020 research and innovation programme under grant agreement No. 871/072.





Goals

- **Determine yields of VCN from candidate** converter materials, including the effects of ND reflectors
- **Determine best converter material for** prototype experiment to be performed at PF1B/ILL
- □ NOTE: ILL PF1B cold spectrum used, some calculations done with generic ILL cold spectrum
- **Results** are in progress and preliminary

	Material	Temperature	Density (g/cm ³)
converter	Solid deuterium	5 K	0.2059
	Liquid parahydrogen	14 K	0.071
	Liquid orthohydrogen	14 K	0.071
reflector	Solid methane	22 K	0.5
	Magnesium-hydride	20 K	0.5 (bulk density)
	Nanodiamond	296 K	0.6 (bulk density)











Solid deuterium cross section

- The SD2 ACE library used in MCNP contains artifacts in energy distribution due to use of (old format) discrete angular and secondary energy distributions.
- Together with the need of better knowledge of SD2 cross section, this motivates the work for improved SD2 cross section (See R. Granada talk).
- □ At present we are using free gas SD2 cross sections at 5.4 K.

Spectrum inside "infinite" sphere with neutrons generated at the center











MCNP geometry of different reflector concepts



Modeling of nanodiamonds in MCNP

- A simple model for SANS as described by Granada [Granada, J. Rolando, J. Ignacio Márquez Damián, and Christian Helman. "Studies on reflector materials for cold neutrons." EPJ Web of Conferences, 231, 04002 (2020)]:
 - Bragg scattering and inelastic scattering are modeled using an ACE file produced with NJOY-Ncrystal
 - SANS is modeled with a piecewise power function that approximates the S(Q) measured by Teshigawara [Teshigawara, M., Y. et al. "Measurement of neutron scattering cross section of nanodiamond with particle diameter of approximately 5 nm in energy range of 0.2 meV to 100 meV." Nucl. Instr. and Meth. A, 929, 113 (2019)]
- The parameters for this function are stored in the ACE file, and a modified version of the Monte Carlo code reads these parameters and sample the outgoing direction.
- The modification was implemented in MCNP 6.2, PHITS 3.21 and OpenMC.

[J. I. Márquez Damián et al., Nuclear Data Development at the ESS, UCANS-web 2020]

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Effect of ND reflector on flux increase



Relative intensity through the tube, solid deuterium converter

λ>20 Å, solid D2



Existing SD2 cross section used

- With nanodiamond around the converter and nanodiamond tube, we can reach 10 times increase in intensity comparing with the bare case.
- The effect is less for a larger converter (green curve)



Angular histogram (orthodeuterium converter and nanodiamond reflector)

Simulation with existing SD2 cross section used and generic ILL cold spectrum



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This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement. No. 871072.

For the prototype experiment we need the intensity at different distances from source, λ > 40 Å





This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement. No. 871072.





Conclusions

□With nanodiamond around the converter and nanodiamond tube, we can reach >10 times increase in intensity comparing with the bare case

The effect is larger for smaller converters (about 3 cm diameter), even though larger converters deliver more VCNs

NEXT STEPS

□For more reliable results, especially above 40 Å, we have identified the need for improved solid deuterium cross sections

Simulations with different ND types will be added in the course of the project

Neutron spectra and count rates of the prototype experiment will be calculated









Second meeting of the Subcommittee on Fundamental Physics of WP3



On the fluorination of nano-diamonds

Marc DUBOIS



How to remove sp² shell and hydrogen?



Flurorination of sp² carbons : 3 effects (crystallinity, SSA, curvature)



How to remove sp² shell and hydrogen?





2 /

Institut de Chimie de Clermont-Ferrand

✓ 1 kg fluorinated nano-diamonds

✓ 2cvcle fluorination (99.9 % F_2 gas, 1 bar, at 450 °C for 12 hours.)



C












¹H MAS NMR (14 kHz)



¹H chemical shift (ppm) vs TMS





7







Fluorination removes the sp² C shell...



Zagrebina et al, J. Phys. Chem. C, 119 (1) (2015) 835–844. Dubois et al, Solid State Nuclear Magnetic Resonance, 40(4) (2011) 144-154. Dubois et al, J. Phys. Chem. C, 113 (24) (2009) 10371–10378





Thanks you for your kind attention

Fluorinated nanodiamonds as unique neutron reflector V. Nesvizhevsky, U. Köster, M. Dubois, N. Batisse, L. Frezet, A.Bosak, L. Gines, O. Williams Carbon, 130 (2018) 799-805

Effect of nanodiamond fluorination on the efficiency of quasispecular reflection of cold neutrons V. V. Nesvizhevsky, M. Dubois, Ph. Gutfreund, E. V. Lychagin, A. Yu. Nezvanov, and K. N. Zhernenkov Phys. Rev. A 97 (2018) 023629





Connecting Russian and European Measures for Large-scale Research Infrastructures Second meeting of the Subcommittee on Fundamental Physics of WP3 8 July 2021

TOWARDS A NEW SCATTERING KERNEL FOR SOLID DEUTERIUM

Rolando Granada





A LETTERS JOURNAL EXPLORING THE FRONTIERS OF PHYSICS

EPL, **86** (2009) 66007 doi: 10.1209/0295-5075/86/66007 June 2009

www.epljournal.org

Neutron scattering kernel for solid deuterium

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The Van Hove scattering function $S(\mathbf{Q}, \omega)$ is directly related to the double-differential cross section:

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{k}{k_0} S(\mathbf{Q}, \omega)$$

 $\begin{array}{ll} {\bf k} \ , \ {\bf k}_{0} & \text{are the scattered and initial neutron wave vectors,} \\ { \hbar \omega } & \text{is the neutron energy loss, and} \\ { \hbar {\bf Q} = \hbar ({\bf k}_{0} - {\bf k}) } \text{ is the momentum transferred to the system} \end{array}$



The scattering law of a molecular system is

$$S(\mathbf{Q},\omega) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dt e^{-i\omega t} \left\langle \sum_{l,l'} \sum_{\nu,\nu'} \overline{a_{l\nu}^* a_{l'\nu'}} \exp\{-i\mathbf{Q}.\mathbf{R}_{l\nu}(0)\} \exp\{i\mathbf{Q}.\mathbf{R}_{l'\nu'}(t)\} \right\rangle$$

where **R** $_{I\nu}$ (t) denotes the position of the atom ν within the molecule I,

$$\mathbf{R}_{l\nu}(t) = \mathbf{a}_l + \mathbf{b}_{\nu}(t) + \mathbf{u}_{\nu}(t)$$

and can be written as the sum of inter $(l \neq l')$ - and intra (l = l')molecular contributions (also referred to as the *outer* and *inner* terms, respectively).



The outer and inner terms, for the intermediate scattering function are

$$\chi(\mathbf{Q},t) = \left\langle \sum_{\substack{l \neq l \\ \nu,\nu'}} \overline{a_{l\nu}^* a_{l\nu'}} \exp\{-i\mathbf{Q}.\mathbf{R}_{l\nu}(0)\} \exp\{i\mathbf{Q}.\mathbf{R}_{l\nu'}(t)\}\right\rangle + \quad \text{"Outer"}$$
$$\left\langle \sum_{\substack{l \\ \nu,\nu'}} \sum_{\substack{\nu,\nu'}} \overline{a_{l\nu}^* a_{l\nu'}} \exp\{-i\mathbf{Q}.\mathbf{R}_{l\nu}(0)\} \exp\{i\mathbf{Q}.\mathbf{R}_{l\nu'}(t)\}\right\rangle \quad \text{"Inner"}$$

Here the brackets denote the average of the time-dependent operators over an equilibrium-distribution function in the full phase space of the scattering system.

In terms of the usual coherent, b_c^{ν} , and incoherent, b_i^{ν} , scattering lengths for nuclei ν ,

$$a_{v} = b_{c}^{v} + 2b_{i}^{v} (\mathbf{S}_{v} \cdot \mathbf{s}) [S_{v} (S_{v} + 1)]^{-1/2}$$



Leaving aside for the moment the consideration of vibrational modes:

$$\chi^{out}(\mathbf{Q},t) = \sum_{l \neq l'} \langle \exp\{-i\mathbf{Q}.\mathbf{a}_{l}(0)\} \exp\{i\mathbf{Q}.\mathbf{a}_{l'}(t)\} \rangle \sum_{v,v'} \langle \overline{b_{lv}^{*}} \exp\{-i\mathbf{Q}.\mathbf{b}_{lv}(0)\} \overline{b_{l'v'}} \exp\{i\mathbf{Q}.\mathbf{b}_{l'v'}(t)\} \rangle$$

$$|\mathbf{d}_{l}(\mathbf{Q},t)| = \sum_{l \neq l'} \langle \mathbf{Q}.\mathbf{d}_{l}(\mathbf{Q},t) | \mathbf{Q}.\mathbf{d}_{l'v'}(t) \rangle \langle \mathbf{Q}.\mathbf{d}_{l'v'}(t) \rangle \langle \mathbf{d}_{l'v'}(t) \rangle \langle \mathbf{d}_{$$

There is no correlation between atoms belonging to different molecules, and then

$$u(\mathbf{Q}) = |\Sigma_v < b_{|v} \exp\{i\mathbf{Q}.b_v(0)\} > |^2 = 4 (b_c)^2 j_0^2 (Qd/2)$$

:. $\chi^{\text{out}}(\mathbf{Q},t) = 4 \ (b_c \)^2 \ j_0^2 \ (Qd/2) \ . \ I_d \ (\mathbf{Q},t)$



$$\chi^{inner}(\mathbf{Q},t) = \sum_{l} \left\langle \exp\{-i\mathbf{Q}\cdot\mathbf{a}_{l}(0)\} \exp\{i\mathbf{Q}\cdot\mathbf{a}_{l}(t)\}\right\rangle \cdot \sum_{JJ''} \sum_{\pi\pi'} A_{\pi\pi'}(J) f_{\pi\pi'}(Q;J,J') e^{-i(\omega_{J}-\omega_{J'})t}$$

$$I_{s}(\mathbf{Q},t) \qquad \qquad \vee (\mathbf{Q},t)$$

$$A_{ee}(J) = c \left\{ 4 b_{c}^{2} + 5/2 b_{l}^{2} \right\} P_{e}(J)$$

$$A_{oo}(J) = (1-c) \left\{ 4 b_{c}^{2} + b_{l}^{2} \right\} P_{o}(J)$$

$$A_{eo}(J) = c \ 3/2 \ b_{l}^{2} P_{e}(J)$$

$$A_{oe}(J) = (1-c) \ 3 \ b_{l}^{2} P_{o}(J)$$

$$f_{\pi=\pi}(Q;J,J) = (2J+1) \ \Sigma_{l \ even}(2l+1) \ C^{2}(JJ';000) \ j_{l}^{2}(Qd/2)$$

$$f_{\pi=\pi}(Q;J,J) = (2J+1) \ \Sigma_{l \ odd}(2l+1) \ C^{2}(JJ';000) \ j_{l}^{2}(Qd/2)$$

$$\therefore \ \chi^{\text{inner}} (\mathbf{Q}, t) = v(\mathbf{Q}, t) \ . \ \mathbf{I}_{s} (\mathbf{Q}, t)$$



$$I_{d}(\mathbf{Q},t) = \sum_{l \neq l'} \langle \exp\{-i\mathbf{Q}.\mathbf{a}_{l}(0)\} \exp\{i\mathbf{Q}.\mathbf{a}_{l'}(t)\} \rangle.$$

$$= \sum_{l,l'} \langle \exp\{-i\mathbf{Q}.\mathbf{a}_{l}(0)\} \exp\{i\mathbf{Q}.\mathbf{a}_{l'}(t)\} \rangle - \sum_{l} \langle \exp\{-i\mathbf{Q}.\mathbf{a}_{l}(0)\} \exp\{i\mathbf{Q}.\mathbf{a}_{l}(t)\} \rangle$$

 $\therefore I_d(\boldsymbol{Q},t) = I(\boldsymbol{Q},t)$ $I_{s}(\boldsymbol{Q},t)$ -

 $\chi(\mathbf{Q},t) = \chi^{\text{out}}(\mathbf{Q},t) + \chi^{\text{inner}}(\mathbf{Q},t)$

$$\chi(\mathbf{Q},t) = 4 \ b_c^2 \ j_0^2 \ (Qr/2) \ \{I \ (\mathbf{Q},t) - I_s \ (\mathbf{Q},t)\} + v(\mathbf{Q},t) \ . \ I_s \ (\mathbf{Q},t)$$

$$\chi(\mathbf{Q},t) = 4 \ b_c^2 \ j_0^2 \ (Qr/2) \ \{I \ (\mathbf{Q},t) - I_s \ (\mathbf{Q},t)\} + v(\mathbf{Q},t) \ . \ I_s \ (\mathbf{Q},t)$$



 $\chi(\mathbf{Q},t) = 4 b_c^2 j_0^2 (Qr/2) \{ I(\mathbf{Q},t) - I_s(\mathbf{Q},t) \} + v(\mathbf{Q},t) \cdot I_s(\mathbf{Q},t) \}$

Now, (must multiply everything by $\chi^{vib}(\mathbf{Q},t)$!)

 $\chi(\mathbf{Q},t) = \chi(\mathbf{Q},0) + \chi(\mathbf{Q},t{\neq}0)$

(elast) (inelast)

but because

 $I_{s}(\mathbf{Q},0) = 1$, $I(\mathbf{Q},0) = |F(\mathbf{Q},0)|^{2}$

 $\chi^{el}(\mathbf{Q},0) = 4 b_c^2 j_0^2 (Qr/2) |F(\mathbf{Q})|^2 + v(\mathbf{Q},0) - u(\mathbf{Q})$

within the Incoherent Approximation:

 $I(\mathbf{Q}, \mathbf{t} \neq \mathbf{0}) \cong I_{s}(\mathbf{Q}, \mathbf{t} \neq \mathbf{0})$

 $\chi^{\text{inel}}(\mathbf{Q},t) = v(\mathbf{Q},t)$. $\mathbf{I}_{\text{s}} \; (\mathbf{Q},t)$



The old calculations were performed according to:

$$\chi^{el}(\mathbf{Q},0) = 4 b_c^2 j_0^2 (Qr/2) |F(\mathbf{Q})|^2 \chi^{vib}(\mathbf{Q},0)$$
 Convent.EI.Coh
+ 2 (1+ α) $b_i^2 \chi^{vib}(\mathbf{Q},0)$ Total EI.Incoh

 $(\alpha = \frac{1}{4} \text{ for } 0 - D2, -\frac{1}{2} \text{ for } p - D2, 0 \text{ for } n - D2)$

 $\chi^{\text{inel}}(\mathbf{Q},t) = \mathbf{v}(\mathbf{Q},t) \cdot \mathbf{I}_{s}(\mathbf{Q},t) \cdot \chi^{\text{vib}}(\mathbf{Q},t)$ Total Inelast.







Whereas the complete elastic components are,

$\chi^{\mathrm{el}}(\mathbf{Q},0) =$	4 $b_c^2 j_0^2$ (Qr/2) $ F(Q) ^2 \chi^{vib}(Q,0)$	Convent.EI.Coh
	+ 2 b _i ² χ ^{vib} (Q ,0)	Convent.El.Incoh
	+ 2 b _i ² α j ₀ (Qd) χ ^{vib} (Q,0)	Spin Correlation
	+ 2 b_c^2 {1+ j_0 (Qd)- 2 j_0^2 (Qd/2)} χ^{vib} (Q,0)	∆ structure

 $(\alpha = \frac{1}{4} \text{ for } o-D2, -\frac{1}{2} \text{ for } p-D2)$



However, a better description of the lattice inelastic component can be attained by preserving the correct expressions for the one-phonon coherent and incoherent terms in the phonon expansion of $I(\mathbf{Q}, t \neq 0)$ and $I_s(\mathbf{Q}, t \neq 0)$ respectively, and applying the incoherent approximation only for the higher phonon terms (≥ 2).

This is the one-phonon correction to the Incoherent Approximation.



Phonon Dispersion Relations (o-D2)



Angular dependence Elas.Incoh.



Phonon Density of states







CONCLUSIONS

- A new scattering kernel to describe the interaction of slow neutrons with solid Deuterium is being developed.
- The main characteristics of this molecular solid are contained in the formalism, including dynamical aspects related to:
 - the lattice's density of states,
 - the Young-Koppel quantum treatment of the rotational motion,
 - the exact treatment of the one-phonon IA in the inelastic term,
 - the internal molecular vibration.
- The elastic processes involving coherent and incoherent contributions are also fully described, as well as the spin-correlation effects caused by the coupling of intrinsic and rotational angular momenta.



CONCLUSIONS

The validation of a new scattering kernel and its generated cross section data demands the comparison with experimental data.

Total cross section measurements are an appropriate set of data to check the theoretical predictions.

The preparation of solid o-Deuterium samples with good repetitive characteristics has to be well understood.



THANKS FOR YOUR ATTENTION!











Reactor PIK and VCN Source

Vladimir Voronin, NRC "Kurchatov institute"-PNPI



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Reactor PIK parameters

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Russian and European Measure for Large-scale Research Infrastructures

Reactor PIK status and plans (optimistic)

Power (W)



 We have permission (Dec. 2020) for 10 MW power from State Regulator (Rostehnadzor).

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- Individual systems for reaching 10 MW are launching with a check of their operability now.
- This year we plan to reach power 10MW.
- Shut down until the end of 2022 to switch to another type of fuel elements, change experimental channels....
- End of 2022 to the permission for 100 MW and the start full power program
- 2024 to start user program



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5 "first day" stations











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Status of the "first day" stations

- 1. TNR test neutron reflectometer
- 2. NERO Polarized Neutron Reflectometer (GKSS)
- 3. DPN polarized neutron diffractometer
- 4. Texture Diffractometer TEX-2 (GKSS)
- 5. Test neutron spectrometer
- Commissioned at the December 2020
- First "demonstration" experiments at 100kW power were carried out
- All results coincide with the theoretical predictions
- Program for updates of these stations to the power of 100 MW is developed











Status of the instrumentation program

- Funding confirmed and started last year (2020)
- The design of the project get the positive state expertise decision in September 2020
- Construction of the first phase instruments started in December 2020
- Second phase instruments construction started June 2021
- The third phase will be started in the beginning of 2022



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tov institute"-PNPI W Road map of the national instrumentation CREMLIN PLUS

		2019 2020			2021				2022				2023				2024						
#	#		IV	I	Ш		IV	I	П		IV	I	П		IV	Ι	Ш	III	IV	I	Ш		IV
(0 Reactor PIK commissioning		100 kW								10	ŴŴ				10-10	0MW	N		~100MW			
	1 Project of instruments																						
1	2 Experimental channel																						
	3 HNC HEC-8																						
4	4 UCNS HEC-4																						
	5 CNS HEC-2																						
e	6 CNS HEC-3																						
	7 Neutronguide system																						
		-						Neut	ron sta	tions								-					
1	SESANS																						
2	INAA								Phas	e 1 (4	instrur	nents)											
3	«Нейтрино» (Neutrino)																						
4	D1																						
5	DC-1																						
6	D3																						
7	IN-1																						0
8	IN-2											Phase 2 (7 instruments)											0
9	ИРИНА (IRINA)																						ommis
10	DEDM																						3
11	FISCO																						N.
12	Tenzor																						<u>N</u> .
13	Мембрана – 2 (Membrane – 2)																						2
14	IN-3																						Į Į.
15	IN-4																						ioning
16	SONATA		I	I	I		I		l														
17			Nuclear physics and particle physics									Phase 3 (9 instruments)											
18			Large scale																				
19			Spectoscopy																				
20			S	pecros	lopy	I	I		I														
	(neutron beta decay)																						

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Hall of experimental channels





N♀	Beam position	Instrument
1	HEC-2	CNS
2	HEC-3	CNS
3	HEC-4	ИУХН (UCN source)
4	HEC-5-5'	ИРИНА (IRINA)
5	HEC-6	
6	HEC-6'	D1
7	HEC-8 (HNS)	DC1
8	HEC-9-1	D3
9	HEC-9-2	ДПН (PolDi)
10	HEC-10	IN-1
11		NEUTRINO

All channels are occupied by the instruments and CN sources from National instrumentational program with the commissioning in 2024-2025. **Only HEC-6 is empty.**

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Problems and questions

Is it a prototype or full scale installation with beam positions and infrastructure?

1. We need a conceptual design of the VCN project with the calculation of radiation shield, weight,

sizes, infrastructure requirements (electric power, air, helium,....). What neutron do you need?

What flux?

Technical

1. Insufficient space for equipment

Administrative

 National instrumental program is almost fixed. The changes are extremely difficult

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Conclusion

Feasibility of VCN source at PIK reactor depends on answer the question:

- What type of source we want to build - prototype or **full scale installation** with beam positions and infrastructure?

To involve PNPI to VCN source project is very easy. Give us:

- A conceptual design of the VCN project and experimental stations which will use the VCN
- Scientific background

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Se NRC "Kurchatov institute"-PNPI W Thank You for attention



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